

**Research Article**

# **BULK VISCOUS BIANCHI-I UNIVERSE WITH DECAYING VACUUM ENERGY**

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## **ABSTRACT**

In the present article, we have investigated Bianchi type-I cosmological model with varying cosmological constant in the presence of bulk viscous fluid by taking the condition  $\zeta = \zeta_0 \theta$ . Einstein's field equations

$$\Lambda = \beta \frac{\ddot{R}}{R} + \frac{1}{R^2}$$

have been solved by assuming the decay law (Overduin and Cooperstock, 1998). Physical and geometrical properties of the model have also been discussed.

**Key Words:** *Bianchi I Space-Time, Bulk Viscous Fluid, Cosmological Constant*

## **INTRODUCTION**

The universe at large scale is homogeneous and isotropic and the accelerating phase of universe (Gasperini, 2003). It is well known that the exact solutions of general theory of relativity for homogeneous space times belongs to Bianchi types. Taking into account dissipative process due to viscosity, the nature of cosmological solutions for homogeneous Bianchi type-I model was investigated by Belinski and Khalatnikov, (1975). They showed that the viscosity can not remove the cosmological singularity but results in a qualitatively new behaviour of the solutions near singularity. They found the remarkable property that during the time of the big bang matter is created by the gravitational field. Bianchi type-I model with bulk viscosity a power function of energy density when the universe is filled with stiff matter were studied by Banerjee,(1985); Huang, (1990). The effect of bulk viscosity with a time varying bulk viscous coefficient, on the evolution of isotropic FRW models in the context of open thermodynamics system was studied by Desikan, (1997).

Models with dynamic cosmological term  $\Lambda(t)$  are becoming popular as they solve the cosmological constant problem in a natural way. There is a significant observational evidence for the detection of Einstein's cosmological constant  $\Lambda$  or a component of material content of the universe that varies slowly with the time and space and so acts like  $\Lambda$ . Some of the recent discussions on the cosmological constant "problem" and consequence on cosmology with a time-varying cosmological constant have been studied by Dolgov, (1993,1997); Sahani and Starobinsky,(2000); Padmnabhan,(2003); Vishwakarma,(1999,2000); Pradhan *et al.*, (2001,2002).

Bulk viscosity is associated with GUT phase transition and string creation. The model studied by Murphy,(1973) has an interesting feature that the big bang type of singularity of infinite space-time curvature does not possess a finite past. However, the relationship assumed by Murphy between the viscosity coefficient and the matter density is not acceptable at large density. The effect of bulk viscosity on cosmological evolution has been investigated by number of authors in the frame work of general relativity (Johri and Sudershan,1988; Maartens, 1995; Zimdahl, 1996). This motivates the study of cosmological bulk viscous fluid model.

In this paper, we have investigated Bianchi type-I cosmological model with varying cosmological

constant in the presence of bulk viscous fluid by taking the condition  $\Lambda = \beta \frac{\ddot{R}}{R} + \frac{1}{R^2}$  model and field

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equation are discussed in section 2. The solution of field equation is given in section 3. Conclusion are summarised in last section.

### METRIC AND FIELD EQUATIONS

We consider the Bianchi type-I space-time represented by the line-element

$$ds^2 = -dt^2 + A^2 dx^2 + B^2 dy^2 + C^2 dz^2 \quad \dots (1)$$

where A, B and C are functions of t only.

We assume the cosmic matter consisting of viscous fluid represented by the energy-momentum tensor

$$T_{ij} = (\rho + \bar{p}) v_i v_j + \bar{p} g_{ij} \quad \dots (2)$$

where  $\bar{p}$  is the effective pressure given by

$$\bar{p} = p - \zeta v_i^i \quad \dots (3)$$

We consider the linear equation of state

$$p = w\rho, \quad 0 \leq w \leq 1 \quad \dots (4)$$

where  $\rho$  and  $p$  are energy density and isotropic pressure respectively and  $v^i$  the four velocity vector of the fluid satisfying  $v_i v^i = -1$ . The Einstein's field equations (in gravitational units  $8\pi G = c = 1$ ) with time-dependent cosmological term  $\Lambda(t)$  are

$$R_i^j - \frac{1}{2} R g_i^j = -T_i^j + \Lambda g_i^j \quad \dots (5)$$

For the line-element (1), the field equations (5) in comoving system of coordinates lead to

$$\frac{\ddot{B}}{B} + \frac{\ddot{C}}{C} + \frac{\dot{B}\dot{C}}{BC} = \Lambda - \bar{p} \quad \dots (6)$$

$$\frac{\ddot{C}}{C} + \frac{\ddot{A}}{A} + \frac{\dot{C}\dot{A}}{CA} = \Lambda - \bar{p} \quad \dots (7)$$

$$\frac{\ddot{A}}{A} + \frac{\ddot{B}}{B} + \frac{\dot{A}\dot{B}}{AB} = \Lambda - \bar{p} \quad \dots (8)$$

$$\frac{\dot{A}\dot{B}}{AB} + \frac{\dot{B}\dot{C}}{BC} + \frac{\dot{A}\dot{C}}{AC} = \Lambda + p \quad \dots (9)$$

The vanishing divergence of Einstein tensor gives rise to

$$\dot{\rho} + (\rho + \bar{p}) \left( \frac{\dot{A}}{A} + \frac{\dot{B}}{B} + \frac{\dot{C}}{C} \right) + \dot{\Lambda} = 0 \quad \dots (10)$$

We define average scale factor R for Bianchi I universe as

$$R^3 = ABC \quad \dots (11)$$

In analogy with FRW universe, we define generalized Hubble parameter H and deceleration parameter q as

$$H = \frac{\dot{R}}{R} = \frac{1}{3} (H_1 + H_2 + H_3) \quad \dots (12)$$

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and  $q = -\frac{\ddot{R}}{RH^2}$  ... (13)

where  $H_1 = \dot{A}/A, H_2 = \dot{B}/B, H_3 = \dot{C}/C$  are directional Hubble factors along x, y and z directions respectively.

We introduce volume expansion  $\theta$  and shear scalar  $\sigma$  for the Bianchi I metric as

$$\theta = v_{ji}^i \quad \dots (14)$$

$$\sigma^2 = \frac{1}{2} \sigma_{ij} \sigma^{ij} \quad \dots (15)$$

For the metric (1), we have

$$\theta = 3\dot{R}/R \quad \dots (16)$$

$$\sigma = K/R^3 \quad \dots (17)$$

where  $3K^2 = k_1^2 + k_2^2 + k_1 k_2$  and  $k_1, k_2$  are integration constant.

Equations (6) – (9) can be expressed in terms of H,  $\sigma$  and q as

$$p - \zeta\theta - \Lambda = (2q - 1)H^2 - \sigma^2 \quad \dots (18)$$

$$\rho + \Lambda = 3H^2 - \sigma^2 \quad \dots (19)$$

## SOLUTION OF THE FIELD EQUATIONS

From equations (4), (18) and (19) we obtain

$$\frac{(1-w)\rho}{2} + \Lambda = \frac{\ddot{R}}{R} + \frac{2\dot{R}^2}{R^2} - \frac{3}{2} \frac{\zeta\dot{R}}{R} \quad \dots (20)$$

Thus, we have one equation with three unknowns R,  $\rho$ ,  $\Lambda$  and  $\zeta$ . We require three more conditions to close the system. We assume that  $p = \rho$  (stiff fluid) i.e.  $w = 1$ . We now consider the decay law for  $\Lambda$  as

$$\Lambda = \beta \frac{\ddot{R}}{R} + \frac{1}{R^2} \quad \dots (21)$$

and bulk viscosity is taken as

$$\zeta = \zeta_0 \theta \quad \dots (22)$$

where  $\beta$  and  $\zeta_0$  are constants.

Therefore, from (20), (21) and (22), we have

$$(1-\beta) \frac{\ddot{R}}{R} + \left(2 - \frac{3}{2} \zeta_0\right) \frac{\dot{R}^2}{R^2} = \frac{1}{R^2} \quad \dots (23)$$

For  $\beta = 1$ , integrating (23), we get

$$R = at + b$$

where  $a = 1/\sqrt{2-9\zeta_0/2}$  and b is an integration constant.

For this solution metric (1) assumes the following form

$$ds^2 = -dt^2 + (at + b)^2 \left[ e^{-\frac{2k_1+k_2}{3a(at+b)^2}} dx^2 + e^{\frac{k_1-k_2}{3a(at+b)^2}} dy^2 + e^{\frac{k_1+2k_2}{3a(at+b)^2}} dz^2 \right]$$

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### DISCUSSION

Matter density  $\rho$  and cosmological term  $\Lambda$  are given by

$$\rho = \frac{3a^2 - 1}{(at + b)^2} - \frac{k^2}{(at + b)^6}$$

$$\Lambda = \frac{1}{(at + b)^2}$$

Expressions for average scale factor  $R$  expansion scalar  $\theta$  shear scalar  $\sigma$ , deceleration parameters  $q$ , Hubble parameter  $H$  and bulk viscosity  $\zeta$  are given by

$$R = at + b$$

$$\theta = \frac{3a}{at + b}$$

$$\sigma = \frac{k}{(at + b)^3}$$

$$q = 0$$

$$H = \frac{a}{at + b}$$

$$\zeta = \frac{3a \zeta_0}{at + b}$$

We observe that model has initial singularity at  $t = -b/a$ . The model starts with a big bang from its initial singularity. At the initial singularity,  $\rho$ ,  $\Lambda$ ,  $\theta$ ,  $\sigma$ ,  $\zeta$ ,  $H$  all are infinite and at late times they become zero. In this model we have  $q = 0$  this indicates that expansion rate of our model is constant. We also observe that  $\sigma / \theta \rightarrow 0$  as  $t \rightarrow \infty$  therefore model approaches isotropy at late times.

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